

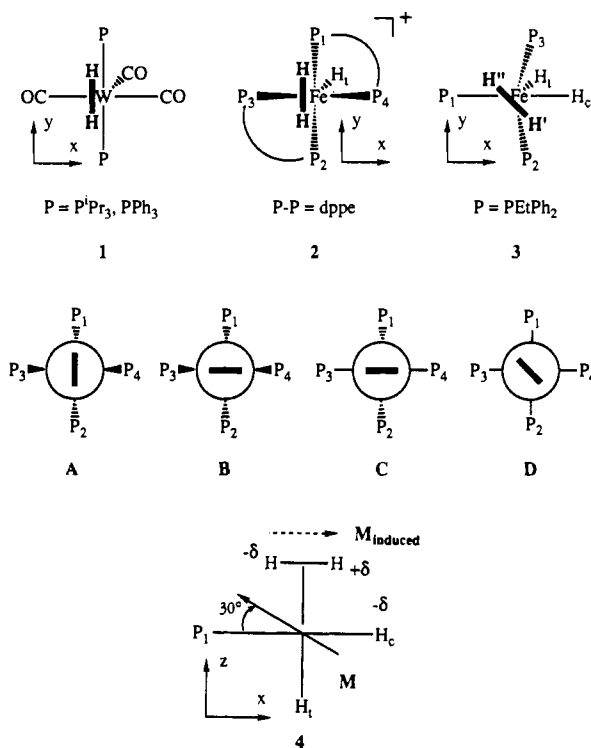
Influence of a Cis Hydride on a Coordinated H₂ Ligand. Ab Initio Calculations

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Numerous transition metal complexes of molecular H₂¹ have been characterized by various spectroscopic means, but determination of accurate H atom positions by neutron diffraction has only been possible in three cases.^{2–4a} In the crystal, the sideways-bonded H₂ ligand is locked into a preferred orientation, but low-temperature INS measurements show rotation of H₂ around the metal–H₂ axis with a barrier less than 2.5 kcal/mol.^{4a,5} The orientations of H₂ are different for the three d⁶ ML₅(H₂) complexes (1–3). H₂ is eclipsed with the P–M–P axis in 1 and



2 and staggered with respect to the four cis ligands in 3. This unusual staggered orientation for an η^2 ligand bonded to a d⁶ ML₅ fragment has been discussed on the basis of EHT calculations.^{4a} Theoretical studies^{4–7} have suggested that two

factors determine the conformation of H₂, i.e., the back-donation from a metal d orbital into σ^*_{HH} (4) and the “cis effect”. Optimal back-donation is obtained when the highest occupied d level with orbital extension toward H₂ overlaps with σ^*_{HH} . The cis effect is an additional two-electron interaction between σ_{Fe-H} and σ^*_{HH} (5), which stabilizes the conformation where H–H eclipses Fe–H. This creates a nascent bond between the closest nonbonded H centers. In the absence of a cis hydride (1^{6,7} and 2^{4a}), back-donation governs the H₂ conformation and the optimal conformation is one that eclipses H–H with the P–M–P axis. In 3, the back-donation puts H₂ in the P1–Fe–P2 plane while the cis effect favors eclipsing of H–H and Fe–H bonds, i.e., in the orthogonal direction. The experimental structure is a consequence of balancing the two effects, and H₂ sits in an intermediate position.^{4a} While back-donation is a well-known metal–ligand interaction, this is not the case for the cis effect,⁴ which may have an important role in the H exchange process.

In order to have a better understanding of the cis effect, we have used ab initio calculations⁸ to study the rotation barrier of H₂ in complex 2, where the cis effect is expected to be absent, and in complex 3, where it should play a role. The crystal structure geometry is used,^{8a} after replacing phosphine ligands with PH₃, unless mentioned otherwise. It is important not to optimize the geometry of the metal fragment in 3, since the ligands are bonded to the metal in a highly distorted octahedral field due to the large bulk of the PEtPh₂ ligand. The rotation of the weakly bonded H₂ ligand is unlikely to modify significantly the geometry of the metallic fragment. We limit ourselves to a section of the potential energy surface in which H₂, at the experimental distance from the Fe center, is allowed to rotate about the Fe–H₂ direction. Our results for 2 and 3 are shown in Figures 1 and 2, respectively.

For H₁(PH₃)₄Fe(H₂)⁺, conformation A is calculated to be the optimal structure. Rotating the ligand by 90° (B) without relaxing the structure of the metal fragment raises the energy by 3.1 kcal/

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- (8) (a) The structure of the H₁(PH₃)₄Fe fragment in 2 has been slightly modified (C_{2v} symmetry) with respect to that in the experimental H(dppe)₂Fe(H₂)⁺ complex. (b) An effective core potential (ECP) was used for Fe and P. The Durand–Barthelat ECP was used for Fe^{11a} with a (3s3p6d) Gaussian basis set contracted into (21/21/411) for the valence shell. Goddard’s 6 Gaussian basis set was used for the d shell.^{12a} The Stevens–Basch ECP was used for P^{11b} with a (4s4p) Gaussian basis set contracted into (31/31) for the valence shell. The H atoms were described with the Huzinaga (4s) Gaussian basis set^{12b} contracted into (31) for H bonded to Fe and into (4) elsewhere. For the H atoms bonded to Fe, one p polarization function ($\zeta = 1.0$) was added. (c) The MCSCF calculations (CASSCF level) used a five-orbital active space for complex 2 (σ_{Fe-H} , σ^*_{Fe-H} , σ_{H-H} , σ^*_{H-H} plus one metal d orbital allowing for back-donation) and a nine-orbital active space for the complex 3 (σ_{Fe-H_1} , σ_{Fe-H_2} , $\sigma^*_{Fe-H_1}$, $\sigma^*_{Fe-H_2}$, σ_{H-H} , σ^*_{H-H} plus three metal d orbitals). We also included the correlation in the metal–hydride bonds, which has an important contribution in a second-order perturbational calculation on the SCF ground state. For complex 3, this small active space is not variationally stable. We had to take into account “intruder states” describing radial correlation by introducing the 3d_{xy} and 4d_{xy} orbitals into the active space. Test calculations adding other d orbitals on the metal or Fe–P bond orbitals in the active space did not change the results of the calculations. (d) A second-order perturbation was carried out on the CASSCF function developed on the five-orbital active space^{9c} following the Möller–Plesset partitioning.¹³ (e) This work is based, in part, on results from the MOTTEC package. Calculations were performed using the HONDO8 package (M. Dupuis, IBM Corp., Center for Scientific & Engineering Computations, Department 48B/428, Neighborhood Rd., Kingston, NY 12401) or a modified version of the MELDF system of programs (E. R. Davidson and D. Feller, QCPE Program No. 580, Quantum Chemistry Program Exchange, Department of Chemistry, Indiana University, Bloomington, IN 47405).

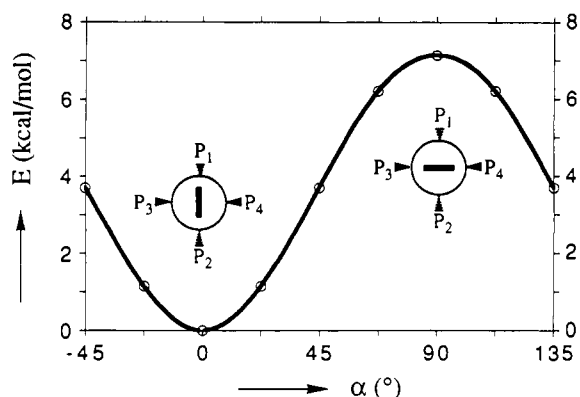


Figure 1. Potential energy (MCSCF) surface (kcal/mol) for rotation of H_2 in **2** as a function of α ($\alpha = 0^\circ$; H-H and P1-Fe-P2 aligned). P1-Fe-P2 = 166.6° and P3-Fe-P4 = 172.1° .³

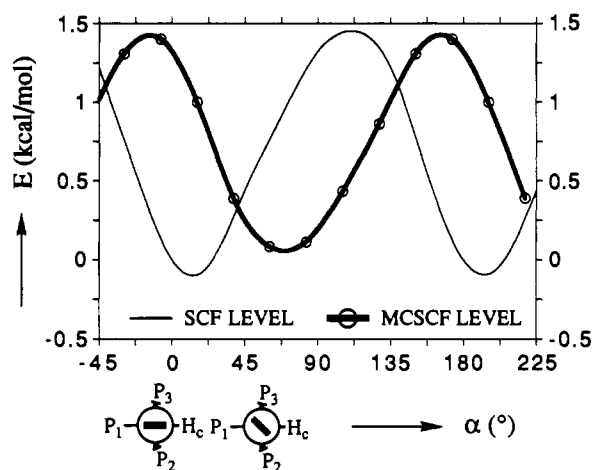


Figure 2. Potential energy surface (kcal/mol) for rotation of H_2 in **3** as a function of α ($\alpha = 0^\circ$; H-H and Fe-H aligned).

mol (SCF) and 7.3 kcal/mol (MCSCF, Figure 1).^{8b} The height of the rotational barrier is 5.9 kcal/mol with CI calculations.^{8d} The origin of the rotational barrier is the diminished back-donation in **B** (the d orbital points away from P3 and P4 and thus away from H_2). It was shown by INS measurements that the rotational barrier can be described mainly by a 2-fold potential of the type $V_2(1 - \cos 2\alpha)$ (where α is the rotational angle for the rotation of H_2 about the Fe- H_2 direction; see Figure 1).^{5b} The height of the rotational barrier can be diminished significantly while its 2-fold shape is maintained, by slightly moving the two phosphines P3 and P4 away from the H_2 ligand (H_c -Fe-P3(4)) = 90° (**C**) in place of 94°). The rotational barrier is then lowered to 1.9 kcal/mol (MCSCF/CI), a value close to the experimental barrier (2.4 kcal/mol).^{5b} Moving the four phosphines as in **D** (H_c -Fe-P = 90°) would also lower the rotational barrier by the same amount, but the barrier would have a dominating $V_4(1 - \cos 4\alpha)$ shape, in contradiction to the INS measurements. These results confirm that the orientation of H_2 is controlled by back-donation. As in **1**,^{5a} it appears that the steric bulk of the phosphine ligand plays no role in the height of the H_2 rotational barrier.

At the SCF level, the best conformation (Figure 2) for **3** corresponds to close alignment of H_2 and the Fe- H_c bond ($\alpha = 12^\circ$). The highest point of the rotational barrier (1.6 kcal/mol) is reached at $\alpha = 110^\circ$, where H_2 almost aligns with the transoid P2-Fe-P3 direction. The inclusion of the correlation energy modifies drastically the previous results. However, at the MCSCF

level,^{8c} the preferred conformation is obtained at $\alpha = 64^\circ$, which is in good agreement with the experimental data, and the top of the barrier (1.4 kcal/mol) is reached for $\alpha = 153^\circ$. The reason for the good agreement between MCSCF calculated and experimental values assuming the rigid-rotator hypothesis is that no ligand cis to H_2 is tilted toward the dihydrogen ligand so that all d orbitals are properly set for efficient back-donation. These two calculations present the two competing interactions at work in this complex: (i) the expected back-donation which, in **3**, places H_2 perpendicular to the Fe-H bond; (ii) an electrostatic interaction favoring coplanarity of H_2 and Fe-H. The latter interaction may be important for the H-exchange process at work in this complex⁴ and in other H_2/H or polyhydride complexes.¹ The electrostatic interaction, well represented at the SCF level, appears via a dipole/induced-dipole interaction. The large dipole moment (4.6 D) of the metallic fragment $H_c(PH_3)_3H_cFe$ lies essentially in the H_c -Fe- H_c plane and makes an angle of 30° with the H_c -Fe-P1 direction (**4**). This induces a dipole moment in the reverse direction on H_2 . The Mulliken analysis shows that the hydrogen center H' closer to H_c is positively charged (0.79) while that on the one H'' is more negatively charged (1.01). The best orientation of H_2 is thus obtained when the two dipoles are coplanar ($\alpha = 0^\circ$). Within EHT methodology, the cis effect results in a weak bonding between the hydride and the closest hydrogen center of H_2 . The Mulliken overlap population of the SCF wave function shows no evidence of a positive bond order.⁹ At the MCSCF level where the back-donation is better calculated,¹⁰ a more accurate balance between the two effects is obtained and the calculated preferred orientation of H_2 is closer to the experimental conformation. This shows the need for high-level correlation calculations for a proper description of the competing factors at work in polyhydride complexes.

Note Added in Proof: After acceptance of our manuscript a theoretical study of the exchange between molecular hydrogen and hydride ligands in *cis*-[Fe(PR₃)₄H(H₂)]⁺ has been published. Masera, F.; Duran, M.; Lledos, A.; Bertran, J. *J. Am. Chem. Soc.* **1992**, *114*, 2922.

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Supplementary Material Available: Tables of structural details and total energies for the optimal conformations at the MCSCF level of complexes **2** and **3** and the ECP and basis set for Fe (3 pages). Ordering information is given on any current masthead page.

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